

# What's New in String Theory?

Steven S. Gubser

Lots of course, including the following:

- Flux compactifications [*Kachru-Kallosh-Linde-Trivedi*]: a huge enough number to solve cosm. const. problem?
- Light cosmic strings [*Copeland-Myers-Polchinski*]: suggested by flux compactifications, may be observable.
- Light strings in preheating [*Gubser*]: a natural but stringy mechanism of reheating.
- String-inspired modification of CDM [*Gubser-Peebles*]: scalar-mediated forces might alter structure formation.

Mainstream of thought on particle phenomenology is still driven by low-energy SUSY and by GUTs.

## Flux compactifications

**The Good News:** Now we can stabilize all moduli of certain Calabi-Yau constructions, break supersymmetry, and get a small cosmological constant.

**The Bad News:** The constructions are complicated, and small  $\Lambda$  is achieved with a shot-gun approach which may limit predictivity.

Compactification on a  $CY_3$  takes us from 10-dim to 4-dim. The  $CY_3$  has *moduli*:

Complex structure:  $\Omega = dz^1 \wedge dz^2 \wedge dz^3$  (locally)

Kahler/size:  $\rho = \text{axion} + iR^4/g_s\alpha'^2$  (assumed unique)

coupling:  $\tau = \text{axion} + i/g_s$  (unique)

(1)

*A priori* these are all massless. After SUSY breaking, size tends to run to  $\infty$ : decompactification.

1) Well-established *flux superpotential* can stabilize  $\Omega$  and  $\tau$ :

$$W_{\text{tree}} = \int G_3 \wedge \Omega \quad G_3 = F_3 - \tau H_3. \quad (2)$$

Choice of  $G_3$  roughly amounts to wrapping 5-branes and then letting them “dissolve” into flux.

2) Instantons or gaugino condensation on other wrapped branes can stabilize  $\rho$ :

$$W = W_{\text{tree}} + Ae^{ia\rho}, \quad (3)$$

where second term is due to  $\langle \lambda\lambda \rangle \sim e^{-8\pi^2/g_{YM}^2 N}$ , with  $\rho = \frac{\theta}{2\pi} + \frac{4\pi}{g_{YM}^2}$ .

But the result of 1) and 2) is unbroken SUSY and  $\Lambda_0 < 0$ .

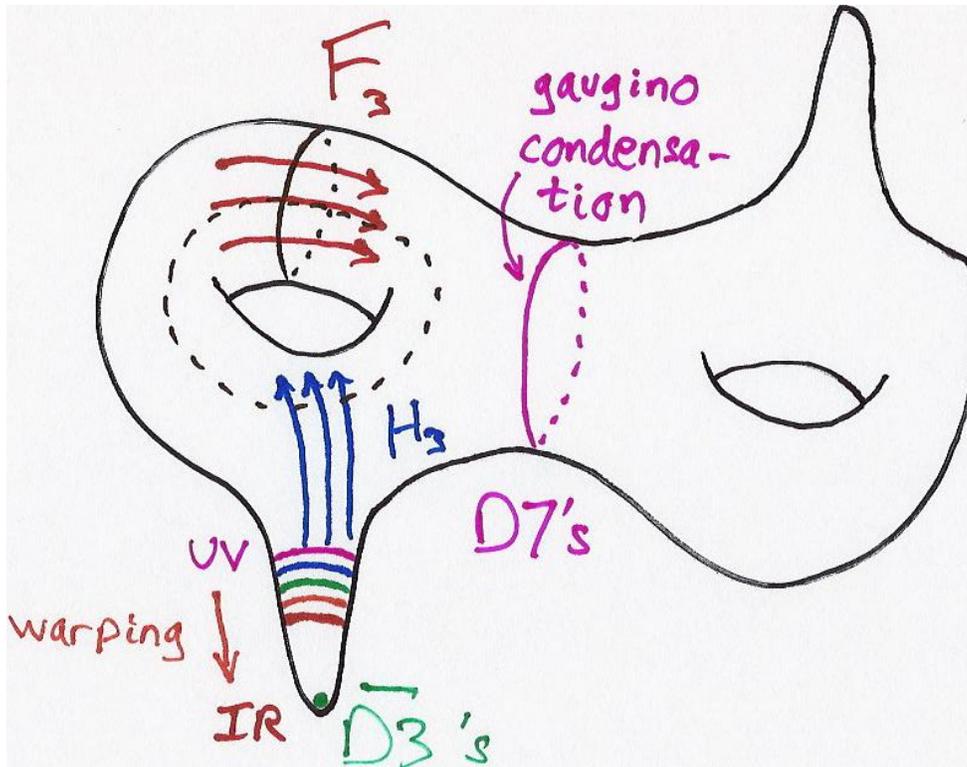
3) Anti-D3-branes can break SUSY and raise  $\Lambda$ :

$$\Lambda = \Lambda_0 + \frac{N_{D3} e^{4A}}{(\text{Im } \rho)^3}, \quad (4)$$

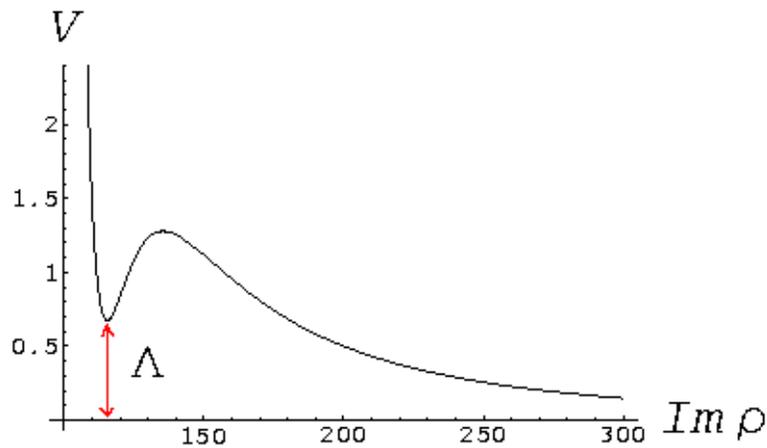
where the *warp factor*  $e^A$  can vary over the  $CY_3$ :

$$ds^2 = e^{2A}(-dt^2 + d\vec{x}^2) + e^{-2A} ds_{CY_3}^2. \quad (5)$$

End picture is complex:



But potential for  $\text{Im } \rho$  is rather nice:



and the key point is that we have *huge but discrete* freedom in choosing  $G_3 \in H^3(\text{CY}_3, \mathbf{Z})$ : we can easily try  $10^{100}$  choices if  $\dim H^3 = 100$ . So small  $\Lambda$  can be arranged. This is the “shotgun” approach.

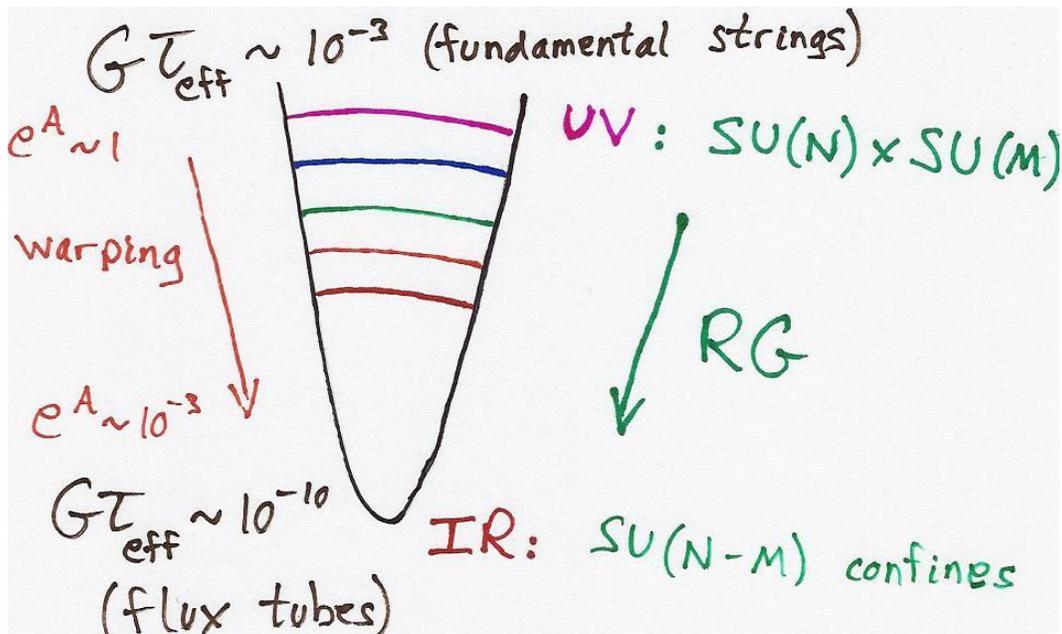
# Light cosmic strings

An old objection to fundamental strings as cosmic strings is

$$\begin{aligned}
 G\tau &\gtrsim 10^{-3} && \text{for fundamental strings} \\
 G\tau &\lesssim 10^{-5} && \text{for cosmic strings,}
 \end{aligned}
 \tag{6}$$

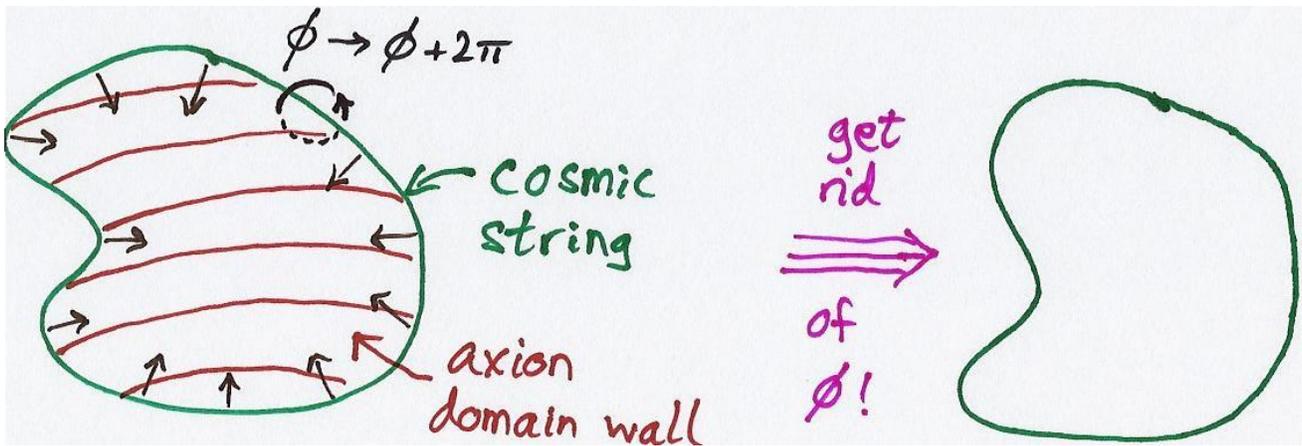
where the second inequality comes from anisotropies in the CMB.

This can be solved in the flux compactification picture thanks to severe warping:

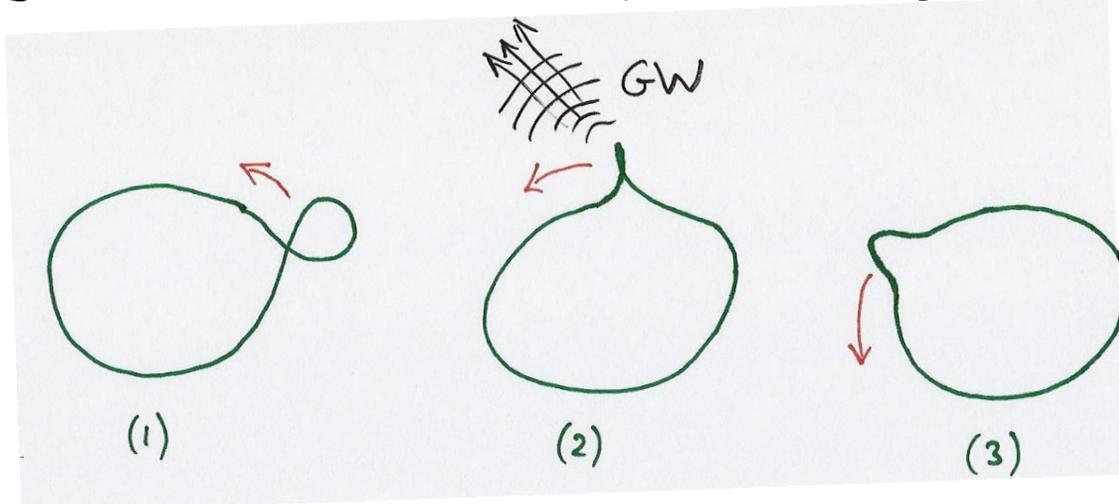


This is not a fine-tuning: dual to a  $SU(N) \times SU(M)$  gauge theory which confines at a scale  $\ll M_{Pl}$ : the cosmic strings in question are basically the flux tubes of this gauge theory.

A further objection based on axionic domain walls is removed by a further feature of flux compactifications: orientifold planes based on a  $Z_2$  which projects away the axion.



A signal for such strings could come from bursts of gravitational radiation produced by cusps:



and LIGO might see this signal for  $G\tau \gtrsim 10^{-11}$ .

## Light strings in preheating

We typically think of creating cosmic strings through the Kibble mechanism (i.e. in a phase transition). What about quantum production of individual vibrational modes?

The on-shell constraint  $L_0|\psi_{\text{phys}}\rangle = 0$  amounts approximately to

$$\ddot{\chi} + \omega^2 \chi = 0, \quad (7)$$

where

$$\omega^2 = k^2 + m^2 \quad m^2 = N/\alpha'. \quad (8)$$

If  $\alpha'$  varies with time, so does  $\omega^2$ , and positive frequency modes for (7) evolve into a mix of positive and negative frequency: *string creation*.

- How might  $\tau = 1/2\pi\alpha'$  vary after inflation?
- Most useful answer is for strings in 4-dim coming from wrapped branes in 10-dim.

Consider D3-branes wrapped on an  $S^2$  whose complexified volume is

$$\varphi \propto \int_{S^2} (B_2 + iJ_2) \quad (\text{a Kahler modulus}). \quad (9)$$

If  $S^2$  shrinks we have light strings in 4-dim. Starting from

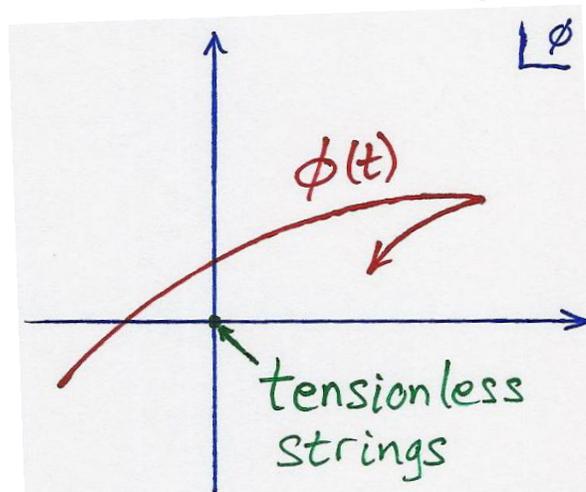
$$S_{D3} = -\tau_{D3} \int d^4\xi \sqrt{G_{\mu\nu} + B_{\mu\nu}} + \dots, \quad (10)$$

with  $\tau_{D3} \sim 1/(\alpha'^2 g_s)$ , obtain

$$\tau_{\text{eff}} \approx M|\varphi| \quad \text{with} \quad M \sim \frac{M_{\text{Pl},4}}{g_s}, \quad (11)$$

where  $\varphi$  has been rescaled to be canonically normalized. (11) is typical for light strings.

String creation after inflation might proceed via a “fly-by” of a point where strings are tensionless:



$$\tau_{\text{eff}} \approx M|\varphi_0 + \dot{\varphi}t| = M\sqrt{|\varphi_0|^2 + (|\dot{\varphi}|t)^2}. \quad (12)$$

To estimate string creation, use a steepest descent trick:

$$\omega^2 \approx m^2 \propto \tau_{\text{eff}} = 0 \quad \text{for } t = -i\mu = -i|\varphi_0|/|\dot{\varphi}|$$

$$|\beta|^2 \approx e^{-\pi\mu\omega(0)}. \quad (13)$$

To get the total number of strings produced, we sum over the Hagedorn spectrum:

$$\frac{dN}{dE} \propto E^\gamma e^{E/T_H}, \quad T_H = \frac{1}{2\pi\sqrt{\alpha' c_\perp/6}}, \quad (14)$$

$$N_{\text{tot}} = \int_0^\infty dE \frac{dN}{dE} |\beta|^2 = \int_0^\infty e^{E/T_H} e^{-\pi\mu\omega(0)}. \quad (15)$$

String production is

$$\begin{aligned} &\text{exponentially suppressed if} && T_H > \frac{1}{\pi\mu} \\ &\text{exponentially enhanced if} && \frac{1}{\pi\mu} > T_H. \end{aligned} \quad (16)$$

The borderline case is when

$$(M\phi_0)^{3/2} \equiv \tau_{\text{min}}^{3/2} = M\dot{\varphi} \sim MM_{\text{Pl},4} m_{\text{inflaton}}. \quad (17)$$

Assuming  $\dot{\varphi} \sim M_{\text{Pl},4} m_{\text{inflaton}}$  is reasonable for oscillations of amplitude  $M_{\text{Pl},4}$  and frequency  $m_{\text{inflaton}}$ .

Further estimating

$$\left. \begin{array}{l} M \sim 10M_{\text{Pl},4} \\ m_{\text{inflaton}} \sim 10^{-5}M_{\text{Pl},4} \end{array} \right\} \implies \sqrt{\tau_{\text{min}}} \lesssim \frac{M_{\text{Pl},4}}{20} \quad (18)$$

in order for strings to be copiously produced right after inflation.

Comparison to *preheating* (coherent production of bosons through parametric resonance driven by inflaton oscillations):

- Causes an even more sudden reheating of the universe.
- Less contrived: (18) is not implausible, and tensionless strings occur on codimension two loci in moduli space.
- Similarly unconstrained by current observations.

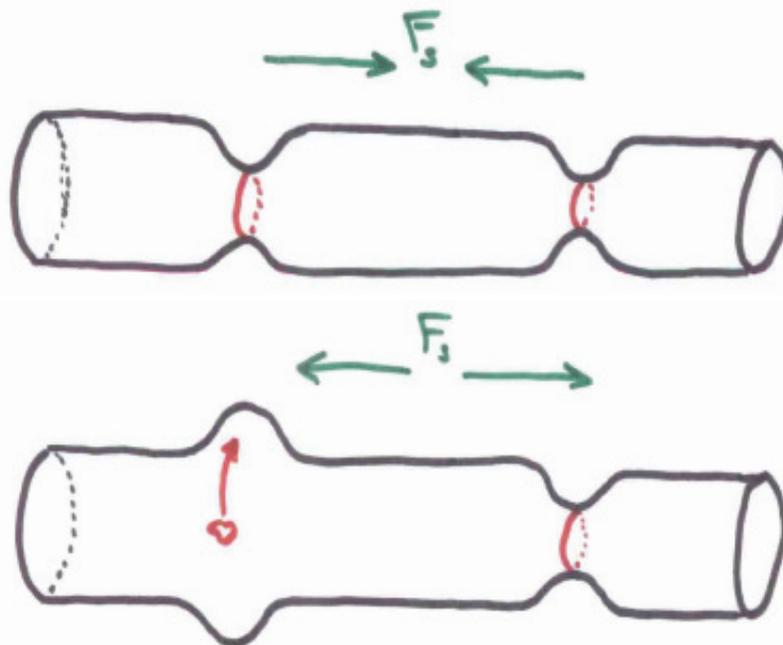
**Speculate** that decay of these strings would cause non-gaussian spectrum of gravitational waves, but I don't have an estimate for LIGO.

## String-inspired modification of CDM

Strings and preheating relied upon rolling scalars at  $10^{13}$  GeV: no conflict with Flux compactifications.

But the next set of ideas hinges on *very light* scalars: if we have to work so hard to get rid of them, why not use them?

Scalar-mediated forces cause like particles to attract and unlike particles to repel. Consider for example strings winding or moving around a compact  $S^1$ :

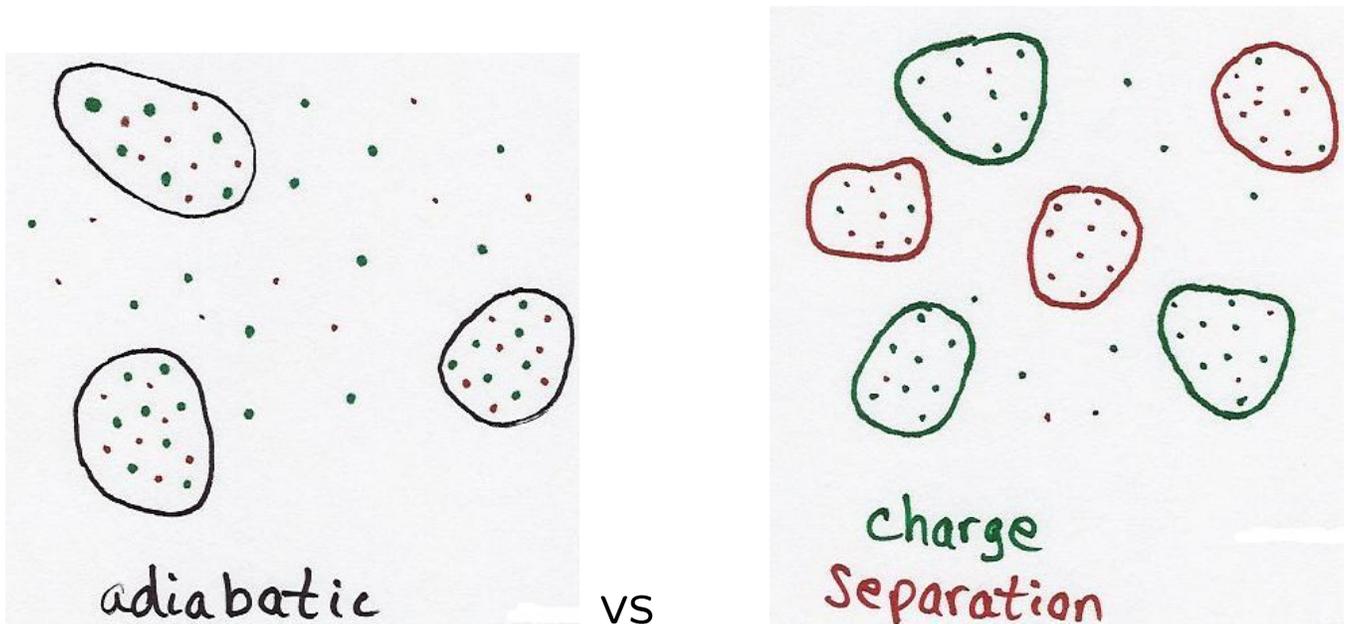


If these forces are comparable to gravity, structure formation might happen differently than in the CDM model.

In general, assume several species of dark matter particles with force law

$$F_{pq} = \beta_{pq} \frac{Gm_p m_q}{r^2} \quad (19)$$

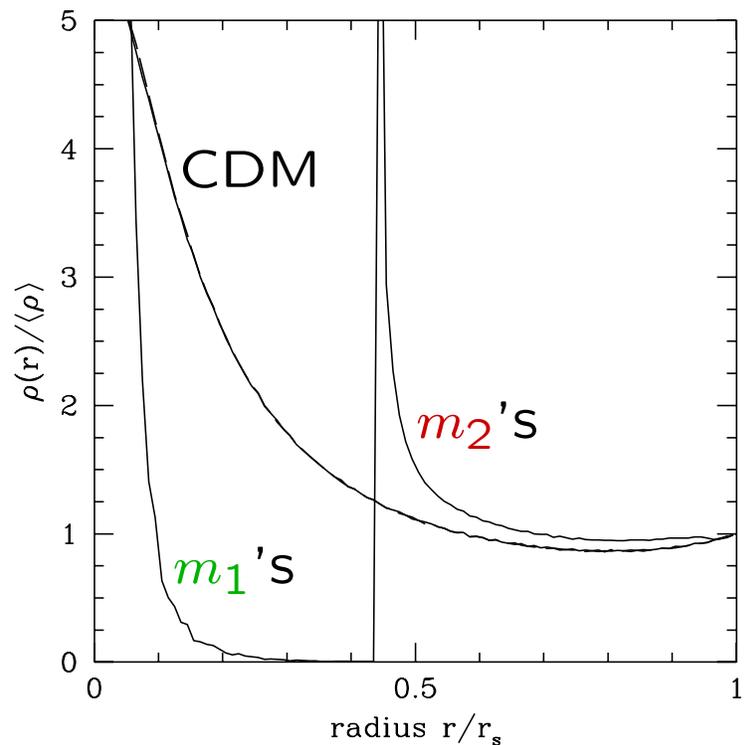
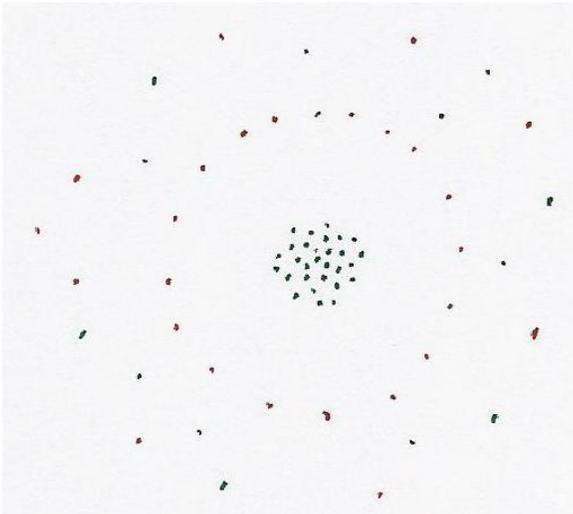
between particles of species  $p$  and  $q$ . Gauge contributions to  $\beta_{pq}$  don't change structure formation (Debye screening), but scalars do:



- Adiabatic mode grows as  $t^{2/3}$ : good match with observations.
- Separation mode grows faster if  $\beta_{12} < 0$ : i.e. if scalar forces are stronger than gravity.
- Naive string theory computation gives  $\beta_{12} = -1$ .
- String theory sits near the border of an observable effect for the linearized regime.

Can allow rather strong scalar forces if they're screened at a scale  $r_s$  comparable to galaxies.

Such forces may ameliorate some problems with CDM. For example, if  $m_1 \ll m_2$ , get an intriguing phenomenon of light vs. heavy halos in bound structures.



Central object is bound by scalar forces, and is much smaller than CDM predicts.

This may be a way of suppressing dwarf galaxies but leaving enough small dense objects to explain recent anomalies in lensing around quasars.

## Summary

- One theme in recent string theory developments has been connection with cosmology.
- Moduli (massless scalars) continue to fascinate, either as heroes or villains.
- Hoping for signals at LIGO, or successes in structure formation.
- With many powerful theoretical tools in hand, still need to focus on the simple and robust features of string theory.

## References

*[Kachru-Kallosh-Linde-Trivedi]* S. Kachru, R. Kallosh, A. Linde and S. P. Trivedi, “De Sitter vacua in string theory,” Phys. Rev. D **68**, 046005 (2003), hep-th/0301240.

*[Copeland-Myers-Polchinski]* E. J. Copeland, R. C. Myers and J. Polchinski, “Cosmic F- and D-strings,” hep-th/0312067.

*[Gubser]* S. S. Gubser, “String creation and cosmology,” hep-th/0312321.

*[Gubser-Peebles]* S. S. Gubser and P. J. E. Peebles, forthcoming.